## String Theory (NS-TP526M) April 22nd 2010

Question 1. $L_{0}-\tilde{L}_{0}$
We have seen that in conformal guage, the Hamiltonian for the closed string is related to the generators of the Virasoro algebra by the identity $\left.H=\left(L_{0}+\tilde{L}\right)_{0}\right) / 2$. Therefore, $L_{0}+\tilde{L}_{0}$ generates time translations on the worldsheet along the time coordinate $\tau$. Similarly, for the closed string, one can express the translations along the spatial coordinate $\sigma$ in terms of $L_{0}-\tilde{L}_{0}$. This follows from the identity

$$
\begin{equation*}
\frac{\mathrm{d} X^{\mu}}{\mathrm{d} \sigma}=\left[-i\left(L_{0}-\tilde{L}_{0}, X^{\mu}\right] .\right. \tag{1}
\end{equation*}
$$

a) Prove this relation, using the identities given in the formularium.

## Question 2. String propagator in light cone guage

Consider the transverse components of the open string operator $X^{i}(\tau, \sigma)$ in light-cone guage. We wish to compute the propagator or two-point correlation function, associated with the transverse modes of an open string. Normal ordering on the oscillator modes is defined by putting all annihilation operators $\alpha_{n}^{i}, n>0$ to the right of the creation operators. Let us extend the definition of normal ordering to the position and momentum operators by: $p^{i} x^{j}:=x^{j} p^{i}$.
a) Show first, using the notation and identities in the formularium for the open string, that

$$
\begin{equation*}
\left[A^{i}(\tau, \sigma), A^{j \dagger}\left(\tau^{\prime}, \sigma^{\prime}\right)\right]=\delta^{i j} G\left(\tau, \tau^{\prime}, \sigma, \sigma^{\prime}\right) \tag{2}
\end{equation*}
$$

where, with the usual worldsheet light-cone coordinates $\sigma^{ \pm}=\tau \pm \sigma$,

$$
\begin{equation*}
G\left(\tau, \tau^{\prime}, \sigma, \sigma^{\prime}\right)=-\frac{1}{4} \log \left[\left(e^{i \sigma^{\prime+}}-e^{i \sigma^{+}}\right)\left(e^{i \sigma^{\prime-}}-e^{i \sigma^{-}}\right)\right]+R\left(\tau, \tau^{\prime}, \sigma, \sigma^{\prime}\right) \tag{3}
\end{equation*}
$$

and $R\left(\tau, \tau^{\prime}, \sigma, \sigma^{\prime}\right)$ denotes terms which are regular in the limit $\tau \rightarrow \tau^{\prime}, \sigma \rightarrow \sigma^{\prime}$.
Find the expression for $G$ and therefore of $R$.
Hint: you may use the identity

$$
\begin{equation*}
\sum_{n=1}^{\infty} \frac{x^{n}}{n}=-\log (1-x) \tag{4}
\end{equation*}
$$

Consequently, show that

$$
\begin{equation*}
X^{i}(\tau, \sigma) X^{j}\left(\tau^{\prime}, \sigma^{\prime}\right)=: X^{i}(\tau, \sigma) X^{j}\left(\tau^{\prime}, \sigma^{\prime}\right):+\delta^{i j} G\left(\tau, \tau^{\prime}, \sigma, \sigma^{\prime}\right) \tag{5}
\end{equation*}
$$

where G is again given by (3), but with a different regular term $R^{\prime}$. Give the expression of $R^{\prime}$ in terms of $R$.

## Question 3. The graviton

In the closed string spectrum in light-cone guage, we have found a physical state

$$
\begin{equation*}
\alpha_{-1}^{i} \tilde{\alpha}_{-1}^{j}|0\rangle, \tag{6}
\end{equation*}
$$

where $i$ and $j$ run over the transverse directions, $i, j=1, \ldots, D-2$.
a) Show that the traceless-symmetric part of this state corresponds to the physical components of a massless graviton in $D$ dimensions. Do this by matching the number of physical degrees of freedom.

## 4. Formularium

### 4.1. Closed string

The oscillator expansion for the closed string in conformal guage reads (in units where the string tension is taken $T=1 / 4 \pi)$,

$$
\begin{equation*}
X^{\mu}(\tau, \sigma)=X_{R}^{\mu}(\tau-\sigma)+X_{L}^{\mu}(\tau+\sigma) \tag{7}
\end{equation*}
$$

where

$$
\begin{align*}
& X_{R}^{\mu}(\tau-\sigma)=\frac{1}{2} x^{\mu}+p^{\mu}(\tau-\sigma)+i \sum_{n \neq 0} \frac{\alpha_{n}^{\mu}}{n} e^{-i n(\tau-\sigma)}, \\
& X_{L}^{\mu}(\tau+\sigma)=\frac{1}{2} x^{\mu}+p^{\mu}(\tau+\sigma)+i \sum_{n \neq 0} \frac{\tilde{\alpha}_{n}^{\mu}}{n} e^{-i n(\tau+\sigma)} . \tag{8}
\end{align*}
$$

The Virasoro generators are normal ordered as

$$
\begin{equation*}
L_{m}=\frac{1}{2}\left(\sum_{n=-\infty}^{+\infty}: \alpha_{m-n}^{\mu} \alpha_{n, \mu}:-a \delta_{m, 0}\right) ; \quad \tilde{L}_{m}=\frac{1}{2}\left(\sum_{n=-\infty}^{+\infty}: \tilde{\alpha}_{m-n}^{\mu} \tilde{\alpha}_{n, \mu}:-a \delta_{m, 0}\right) \tag{9}
\end{equation*}
$$

with $\alpha_{0}^{\mu}=\tilde{\alpha}_{0}^{\mu}=p^{\mu}$. The non-vanishing commutation relations in conformal guage are (we set $\hbar=1$ )

$$
\begin{equation*}
\left[x^{\mu}, p^{\nu}\right]=i \eta^{\mu \nu}, \quad\left[\alpha_{m}^{\mu}, \alpha_{n}^{\nu}\right]=\left[\tilde{\alpha}_{m}^{\mu}, \tilde{\alpha}_{n}^{\nu}\right]=m \delta_{m+n, 0} \eta^{\mu \nu} \tag{10}
\end{equation*}
$$

### 4.2. Open string in light-cone

We can write the oscillator expansion for the open string in light-cone guage as

$$
\begin{equation*}
X^{i}(\tau, \sigma)=x^{i}+p^{i} \tau+A^{i}(\tau, \sigma)+A^{i \dagger}(\tau, \sigma), \tag{11}
\end{equation*}
$$

with

$$
\begin{equation*}
A^{i}(\tau, \sigma)=i \sum_{n=1}^{\infty} \frac{\alpha_{n}^{i}}{n} e^{-i n \tau} \cos (n \sigma), \quad A^{i \dagger}(\tau, \sigma)=-i \sum_{n=1}^{\infty} \frac{\alpha_{-n}^{i}}{n} e^{i n \tau} \cos (n \sigma), \tag{12}
\end{equation*}
$$

with the usual commutation relations

$$
\begin{equation*}
\left[x^{i}, p^{j}\right]=i \delta^{i j}, \quad\left[\alpha_{m}^{i}, \alpha_{n}^{j}\right]=m \delta_{m+n, 0} \delta^{i j} \tag{13}
\end{equation*}
$$

